
Quantum recurrence and matrix Schur functions

Luis Velázquez

Departamento de Matemática Aplicada & IUMA, U Zaragoza

Joint works with:

Reinhard Werner, Albert Werner

Institut für Theoretische Physik, Leibniz Universität Hannover

Jean Bourgain

Institute for Advanced Study, Princeton

Alberto Grünbaum, Jon Wilkening

Department of Mathematics, UC Berkeley

Recurrence and expected return time

A key concept in the study of random processes is the idea of **recurrence**.

A random process is recurrent if it returns to the initial state with probability one. Otherwise it is called transient.

In principle the recurrence of a random process could depend on the initial state, so we should talk about recurrent or transient states for a given random system.

Recurrent states can have different return speeds, which can be measured by the **expected return time**.

Recurrent states can have a finite or infinite expected return time. The last ones are the recurrent states which are closest to the transience.

Random Walks: Pólya recurrence

RW = discrete classical random process \rightsquigarrow discrete-time Markov chain
on a countable state space X

- **Transition matrix**

$$P = (P_{x,y})_{x,y \in X} \quad P_{x,y} \geq 0 \quad \sum_{y \in X} P_{x,y} = 1$$

$P_{x,y}$ = probability of transition $x \rightarrow y$ in one step

- **Return probability in n steps**

$$p_n = p_n^{(x)} = \sum_{x_1, \dots, x_{n-1}} P_{x,x_1} P_{x_1,x_2} \cdots P_{x_{n-1},x} = (P^n)_{x,x}$$

First time

- **return probability in n steps**

$$q_n = q_n^{(x)} = \sum_{x_1, \dots, x_{n-1} \neq x} P_{x,x_1} P_{x_1,x_2} \cdots P_{x_{n-1},x}$$

$$q = \sum_{n=1}^{\infty} q_n$$

Total return probability (*Pólya number*)

x is recurrent $\Leftrightarrow q = 1$

$$\tau = \sum_{n=1}^{\infty} n q_n$$

Expected return time

Random Walks: Pólya recurrence

The study of recurrence is greatly aided by the use of **generating functions**

$$p(z) = \sum_{n=0}^{\infty} p_n z^n \qquad q(z) = \sum_{n=1}^{\infty} q_n z^n$$

A key result relates the first return prob. q_n to the more accessible return prob. p_n

$$q(z) = 1 - \frac{1}{p(z)} \qquad \text{Renewal equation}$$

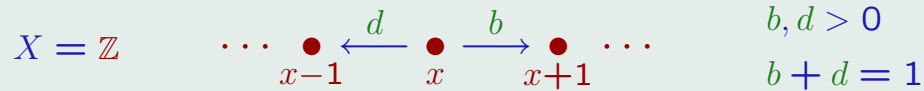
This allows us to rewrite the recurrence notions in terms of p_n instead of q_n

$$q = \sum_{n=1}^{\infty} q_n = q(1) = 1 - \frac{1}{p(1)} \qquad \text{Pólya number}$$

$$x \text{ is recurrent} \Leftrightarrow \sum_{n=0}^{\infty} p_n = \infty$$

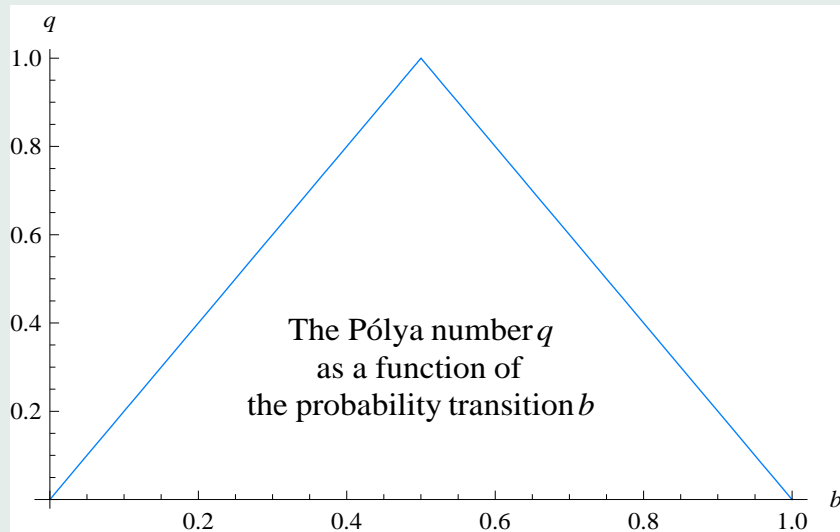
$$\tau = \sum_{n=1}^{\infty} n q_n = q'(1) = \lim_{z \rightarrow 1} \frac{p'(z)}{p(z)^2} \qquad \text{Expected return time}$$

Example: Translation invariant birth-death process



In this case explicit results are available, giving for any state x

$$p(z) = \frac{1}{\sqrt{1 - 4bdz^2}} \quad q = 1 - \sqrt{1 - 4bd}$$



All the states are recurrent if $b = d = \frac{1}{2}$ and transient otherwise

Random Walks: Spectral characterization of recurrence

For **reversible** Markov chains the transition matrix P is **symmetrizable** spectral theory →

For any state $x \in X$ there exists a **spectral measure** m_x on $[-1, 1]$ such that

$$n\text{-th moment } \int_{-1}^1 t^n dm_x(t) = (P^n)_{x,x} \text{ return prob. of } x \text{ in } n \text{ steps}$$

This allows us the use of spectral techniques for RW, but also provides a link to the machinery of orthogonal polynomials on \mathbb{R} .

How is the recurrence of a state x codified in its spectral measure $m = m_x$?

$$x \text{ is recurrent} \Leftrightarrow \int_{-1}^1 \frac{dm(t)}{1-t} = \infty$$
$$\tau = \frac{1}{m(\{1\})} \quad \text{Expected return time}$$

Finite expected return time $\Leftrightarrow 1$ is a mass point

Quantum Walks: Why?

- **What is a QW?** Simplified model of a **quantum system in evolution**
- **Why are QW important?** Simplest setting to study the quantum behaviour:
 - **Superposition:** quantum systems can be simultaneously in different states
 - **Parallelism:** quantum systems may spread much faster than classical ones
 - **Entanglement:** quantum systems arbitrarily far from each other can be linked

Quantum Walks: Why?

- **What is a QW?** Simplified model of a **quantum system in evolution**
- **Why are QW important?** Simplest setting to study the quantum behaviour:
 - **Superposition:** quantum systems can be simultaneously in different states
 - **Parallelism:** quantum systems may spread much faster than classical ones
 - **Entanglement:** quantum systems arbitrarily far from each other can be linked

The exponentially increasing interest in QW rests on the discovery of **astonishing applications to our daily life** of this “**magic**” **quantum behaviour** of nature: **super-phenomena** (photosynthesis extreme efficiency), **quantum computation**, **quantum cryptography**, **quantum teleportation**, ...

Quantum Walks: Why?

- **What is a QW?** Simplified model of a **quantum system in evolution**
- **Why are QW important?** Simplest setting to study the quantum behaviour:
 - **Superposition:** quantum systems can be simultaneously in different states
 - **Parallelism:** quantum systems may spread much faster than classical ones
 - **Entanglement:** quantum systems arbitrarily far from each other can be linked

The exponentially increasing interest in QW rests on the discovery of **astonishing applications to our daily life** of this **“magic” quantum behaviour** of nature: **super-phenomena** (photosynthesis extreme efficiency), **quantum computation**, **quantum cryptography**, **quantum teleportation**, ...

	RW	QW
Randomness	Ignorance	Intrinsic
States	≠ superposition	∃ superposition
Evolution	Markov process	Unitary process
Measurement	Does not disturb evolution	Projection which alter evolution
Methods	Path counting Fourier analysis Spectral self-adjoint - OP line	Path counting Fourier analysis Spectral unitary - OP circle

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

- **States**

If the elements $x \in X$ are all the possible outcomes of a measurement the measurable states form a countable orthonormal basis $\{\phi_x\}_{x \in X}$ of H

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

- **States**

If the elements $x \in X$ are all the possible outcomes of a measurement the measurable states form a countable orthonormal basis $\{\phi_x\}_{x \in X}$ of H

A system can be in a complex **superposition** $\psi = \sum_{x \in X} c_x \phi_x$ and then any measurement of X yields a **random** result with **Prob**(x) = $|c_x|^2$

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

• States

If the elements $x \in X$ are all the possible outcomes of a measurement the measurable states form a countable orthonormal basis $\{\phi_x\}_{x \in X}$ of H

A system can be in a complex **superposition** $\psi = \sum_{x \in X} c_x \phi_x$ and then any measurement of X yields a **random** result with $\text{Prob}(x) = |c_x|^2$

Consequences

► $\|\psi\|^2 = \sum_{x \in X} |c_x|^2 = \text{Prob}(X) = 1 \quad \Rightarrow \quad \psi$ must be unitary!

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

• States

If the elements $x \in X$ are all the possible outcomes of a measurement the measurable states form a countable orthonormal basis $\{\phi_x\}_{x \in X}$ of H

A system can be in a complex **superposition** $\psi = \sum_{x \in X} c_x \phi_x$ and then any measurement of X yields a **random** result with $\text{Prob}(x) = |c_x|^2$

Consequences

- ▶ $\|\psi\|^2 = \sum_{x \in X} |c_x|^2 = \text{Prob}(X) = 1 \quad \Rightarrow \quad \psi$ must be unitary!
- ▶ $\psi \equiv e^{i\zeta} \psi$ same state \Rightarrow State space = complex projective $\text{CP}(\mathcal{H})$

Quantum Walks

QW = discrete quantum evolution \rightsquigarrow discrete-time unitary process on a complex Hilbert space \mathcal{H} with inner product $\langle \cdot | \cdot \rangle$

• States

If the elements $x \in X$ are all the possible outcomes of a measurement the measurable states form a countable orthonormal basis $\{\phi_x\}_{x \in X}$ of H

A system can be in a complex **superposition** $\psi = \sum_{x \in X} c_x \phi_x$ and then any measurement of X yields a **random** result with $\text{Prob}(x) = |c_x|^2$

Consequences

- ▶ $\|\psi\|^2 = \sum_{x \in X} |c_x|^2 = \text{Prob}(X) = 1 \quad \Rightarrow \quad \psi$ must be unitary!
- ▶ $\psi \equiv e^{i\zeta} \psi$ same state \Rightarrow State space = complex projective $\text{CP}(\mathcal{H})$

The non trivial geometry of the state space has measurable consequences!

Quantum Walks

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$

Quantum Walks

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$

▶ The value x is NOT found: Then $\psi \xrightarrow{\text{collapse}}$ projection of ψ onto ϕ_x^\perp

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

- ▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$
- ▶ The value x is NOT found: Then $\psi \xrightarrow{\text{collapse}}$ projection of ψ onto ϕ_x^\perp

$c_x = \langle \phi_x | \psi \rangle$ is called the *amplitude* for the transition $\psi \rightarrow \phi_x$

$$\text{Prob}(\psi \rightarrow \phi_x) = |\langle \phi_x | \psi \rangle|^2$$

Quantum Walks

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

- ▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$
- ▶ The value x is NOT found: Then $\psi \xrightarrow{\text{collapse}}$ projection of ψ onto ϕ_x^\perp

$c_x = \langle \phi_x | \psi \rangle$ is called the *amplitude* for the transition $\psi \rightarrow \phi_x$

$$\text{Prob}(\psi \rightarrow \phi_x) = |\langle \phi_x | \psi \rangle|^2$$

- **Evolution = Unitary operator**

The one-step evolution is given by $\psi \rightarrow U\psi$ where U is a unitary operator on \mathcal{H}

Unitarity ensures conservation of total probability $\|\psi\|^2 = \text{Prob}(X) = 1$

Quantum Walks

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

- ▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$
- ▶ The value x is NOT found: Then $\psi \xrightarrow{\text{collapse}}$ projection of ψ onto ϕ_x^\perp

$c_x = \langle \phi_x | \psi \rangle$ is called the *amplitude* for the transition $\psi \rightarrow \phi_x$

$$\text{Prob}(\psi \rightarrow \phi_x) = |\langle \phi_x | \psi \rangle|^2$$

- **Evolution = Unitary operator**

The one-step evolution is given by $\psi \rightarrow U\psi$ where U is a unitary operator on \mathcal{H}

Unitarity ensures conservation of total probability $\|\psi\|^2 = \text{Prob}(X) = 1$

$\langle \phi_x | U^n \psi \rangle$ is the amplitude for the transition $\psi \rightarrow \phi_x$ in n steps

$$\text{Prob}(\psi \xrightarrow{n \text{ steps}} \phi_x) = |\langle \phi_x | U^n \psi \rangle|^2$$

Quantum Walks

- **Measurement = Orthogonal projector**

If a system is in a state ψ , the measurement of a concrete value x can result in:

- ▶ The value x is found: Then $\psi \xrightarrow{\text{collapse}} \phi_x$
- ▶ The value x is NOT found: Then $\psi \xrightarrow{\text{collapse}}$ projection of ψ onto ϕ_x^\perp

$c_x = \langle \phi_x | \psi \rangle$ is called the *amplitude* for the transition $\psi \rightarrow \phi_x$

$$\text{Prob}(\psi \rightarrow \phi_x) = |\langle \phi_x | \psi \rangle|^2$$

- **Evolution = Unitary operator**

The one-step evolution is given by $\psi \rightarrow U\psi$ where U is a unitary operator on \mathcal{H}

Unitarity ensures conservation of total probability $\|\psi\|^2 = \text{Prob}(X) = 1$

$\langle \phi_x | U^n \psi \rangle$ is the amplitude for the transition $\psi \rightarrow \phi_x$ in n steps

$$\text{Prob}(\psi \xrightarrow{n \text{ steps}} \phi_x) = |\langle \phi_x | U^n \psi \rangle|^2$$

$$\text{Prob} = |\text{amplitude}|^2$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude
in n steps**

$$\mu_n^\psi = \langle \psi | U^n \psi \rangle \quad \rightsquigarrow \quad \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi \xrightarrow{\text{step 2}} U(QU)\psi$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi \xrightarrow{\text{step 2}} U(QU)\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} (QU)^2\psi$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi \xrightarrow{\text{step 2}} U(QU)\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} (QU)^2\psi \xrightarrow{\text{step 3}} \dots$$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q = \text{projector onto } \psi^\perp \equiv \text{collapse for NO measurement of } \psi$

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi \xrightarrow{\text{step 2}} U(QU)\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} (QU)^2\psi \xrightarrow{\text{step 3}} \dots$$

- **First time return amplitude in n steps** $a_n^\psi = \langle \psi | U(QU)^{n-1} \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow[\text{1st time}]{n \text{ steps}} \psi) = |a_n^\psi|^2$

Quantum Walks: Recurrence

Given a unitary step U and an initial state ψ we search for the return probability to ψ

- **Return amplitude in n steps** $\mu_n^\psi = \langle \psi | U^n \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$

However, to avoid overcounting, the total return probability needs the **first time** ones:

$Q =$ projector onto $\psi^\perp \equiv$ collapse for NO measurement of ψ

$$\psi \xrightarrow{\text{step 1}} U\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} QU\psi \xrightarrow{\text{step 2}} U(QU)\psi \xrightarrow[\text{NO } \psi]{\text{measurement}} (QU)^2\psi \xrightarrow{\text{step 3}} \dots$$

- **First time return amplitude in n steps** $a_n^\psi = \langle \psi | U(QU)^{n-1} \psi \rangle \rightsquigarrow \text{Prob}(\psi \xrightarrow[\text{1st time}]{n \text{ steps}} \psi) = |a_n^\psi|^2$

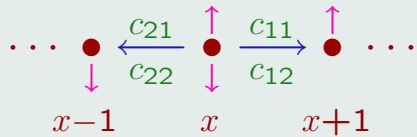
$$R_\psi = \sum_{n=1}^{\infty} |a_n^\psi|^2 \quad \text{Total return probability}$$

$$\psi \text{ is recurrent} \Leftrightarrow R_\psi = 1$$

$$\tau_\psi = \sum_{n=1}^{\infty} n |a_n^\psi|^2 \quad \text{Expected return time}$$

Example: Translation invariant coined walk

Consider a particle with *spin* in an infinite 1D lattice with the one-step amplitudes

$$X = \mathbb{Z} \times \{\uparrow, \downarrow\}$$

$$C = \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix} \text{ unitary coin}$$

Example: Translation invariant coined walk

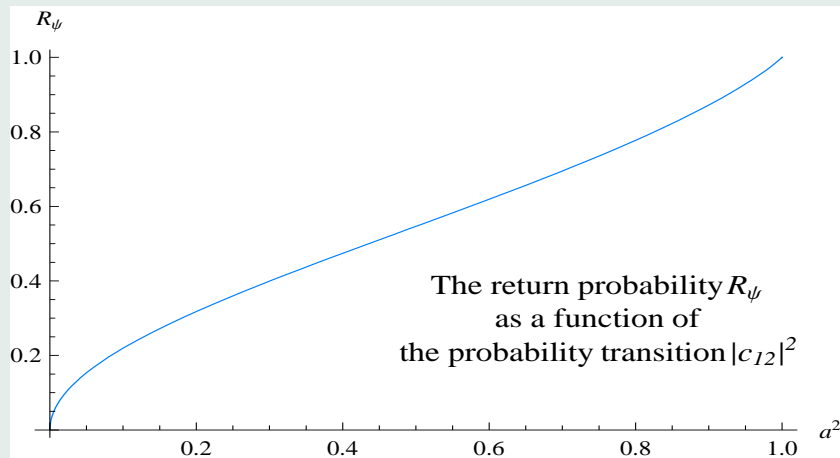
Consider a particle with *spin* in an infinite 1D lattice with the one-step amplitudes

$$X = \mathbb{Z} \times \{\uparrow, \downarrow\} \quad \cdots \quad \begin{array}{c} \bullet \\ \downarrow \\ x-1 \end{array} \quad \begin{array}{c} \xleftarrow{c_{21}} \\ \bullet \\ \downarrow \\ x \end{array} \quad \begin{array}{c} \uparrow \\ \xrightarrow{c_{11}} \\ \bullet \\ \downarrow \\ x+1 \end{array} \quad \cdots$$

$$C = \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix} \text{ unitary coin}$$

Explicit calculations are possible, giving for any basis state $\psi = |x, \uparrow\rangle, |x, \downarrow\rangle$

$$R_\psi = \frac{(1 + 2\rho^2)\rho a + (1 - 4\rho^2) \arcsin a}{\frac{\pi}{2}a^4} \quad \left[\begin{array}{l} a = |c_{12}| \\ \rho = \sqrt{1 - a^2} \end{array} \right]$$



In contrast to the classical case, no recurrent states appear

Quantum Walks: Spectral characterization of recurrence

In contrast to the classical case the transition matrix U is always unitary spectral theory →

For any state $\psi \in \mathcal{H}$ there exists a **spectral measure** μ_ψ on $\mathbb{T} \equiv |z| = 1$ such that

$$n\text{-th moment} \quad \int_{\mathbb{T}} t^n d\mu_\psi(t) = \langle \psi | U^n \psi \rangle = \mu_n^\psi \quad \text{return amplitude to } \psi \text{ in } n \text{ steps}$$

Quantum Walks: Spectral characterization of recurrence

In contrast to the classical case the transition matrix U is always unitary → spectral theory

For any state $\psi \in \mathcal{H}$ there exists a **spectral measure** μ_ψ on $\mathbb{T} \equiv |z| = 1$ such that

$$n\text{-th moment} \quad \int_{\mathbb{T}} t^n d\mu_\psi(t) = \langle \psi | U^n \psi \rangle = \mu_n^\psi \quad \text{return amplitude to } \psi \text{ in } n \text{ steps}$$

This leads to a fruitful link between QW and OP on \mathbb{T} , but also opens the question:

How is the recurrence of a state ψ codified in its spectral measure μ_ψ ?

Quantum Walks: Spectral characterization of recurrence

In contrast to the classical case the transition matrix U is always unitary → spectral theory

For any state $\psi \in \mathcal{H}$ there exists a **spectral measure** μ_ψ on $\mathbb{T} \equiv |z| = 1$ such that

$$n\text{-th moment} \quad \int_{\mathbb{T}} t^n d\mu_\psi(t) = \langle \psi | U^n \psi \rangle = \mu_n^\psi \quad \text{return amplitude to } \psi \text{ in } n \text{ steps}$$

This leads to a fruitful link between QW and OP on \mathbb{T} , but also opens the question:

How is the recurrence of a state ψ codified in its spectral measure μ_ψ ?

The answer comes from the study of the **generating functions**

$$S_\psi(z) = \sum_{n=0}^{\infty} \mu_n^\psi z^n = \int_{\mathbb{T}} \frac{d\mu_\psi(t)}{1 - tz}$$

Stieltjes function of μ_ψ

$$g_\psi(z) = \sum_{n=1}^{\infty} a_n^\psi z^n$$

???

Quantum Walks: Spectral characterization of recurrence

In contrast to the classical case the transition matrix U is always unitary spectral theory →

For any state $\psi \in \mathcal{H}$ there exists a **spectral measure** μ_ψ on $\mathbb{T} \equiv |z| = 1$ such that

$$n\text{-th moment} \quad \int_{\mathbb{T}} t^n d\mu_\psi(t) = \langle \psi | U^n \psi \rangle = \mu_n^\psi \quad \text{return amplitude to } \psi \text{ in } n \text{ steps}$$

This leads to a fruitful link between QW and OP on \mathbb{T} , but also opens the question:

How is the recurrence of a state ψ codified in its spectral measure μ_ψ ?

The answer comes from the study of the **generating functions**

$$S_\psi(z) = \sum_{n=0}^{\infty} \mu_n^\psi z^n = \int_{\mathbb{T}} \frac{d\mu_\psi(t)}{1 - tz} \quad \text{Stieltjes function of } \mu_\psi$$

$$g_\psi(z) = \sum_{n=1}^{\infty} a_n^\psi z^n = z \overline{f_\psi(z)} \quad f_\psi(z) = \text{Schur function of } \mu_\psi!$$

Quantum Walks and Schur functions

A Schur function is a function $f: \mathbb{D} \rightarrow \overline{\mathbb{D}}$ analytic on the unit disk $\mathbb{D} \equiv |z| < 1$

Quantum Walks and Schur functions

A Schur function is a function $f: \mathbb{D} \rightarrow \overline{\mathbb{D}}$ analytic on the unit disk $\mathbb{D} \equiv |z| < 1$

They are in one-to-one correspondence with the measures μ on the unit circle \mathbb{T} via

$$f(z) = \frac{1 - F(z)}{z F(z) + 1} \quad F(z) = \int_{\mathbb{T}} \frac{t + z}{t - z} d\mu(t)$$

Quantum Walks and Schur functions

A Schur function is a function $f: \mathbb{D} \rightarrow \overline{\mathbb{D}}$ analytic on the unit disk $\mathbb{D} \equiv |z| < 1$

They are in one-to-one correspondence with the measures μ on the unit circle \mathbb{T} via

$$f(z) = \frac{1 - F(z)}{z F(z) + 1} \quad F(z) = \int_{\mathbb{T}} \frac{t + z}{t - z} d\mu(t)$$

This class of functions is a classical ingredient in complex analysis and OP on \mathbb{T}

The **NEW** result is the discovery of the **quantum role** of Schur functions:

Schur Taylor coefficients = First time return amplitudes of QW

Quantum Walks and Schur functions

A Schur function is a function $f: \mathbb{D} \rightarrow \overline{\mathbb{D}}$ analytic on the unit disk $\mathbb{D} \equiv |z| < 1$

They are in one-to-one correspondence with the measures μ on the unit circle \mathbb{T} via

$$f(z) = \frac{1 - F(z)}{z F(z) + 1} \quad F(z) = \int_{\mathbb{T}} \frac{t + z}{t - z} d\mu(t)$$

This class of functions is a classical ingredient in complex analysis and OP on \mathbb{T}

The **NEW** result is the discovery of the **quantum role** of Schur functions:

Schur Taylor coefficients = First time return amplitudes of QW

The identification of $g_\psi(z) = z \overline{f_\psi}(z)$ as a Schur function is a consequence of

$$g_\psi(z) = 1 - \frac{1}{S_\psi(z)}$$

**Quantum
renewal equation**

Like the classical one!
But for amplitudes
instead of probabilities

and is **KEY** for the spectral analysis of quantum recurrence

Quantum Walks: Spectral characterization of recurrence

Denoting $\langle f, g \rangle = \int_0^{2\pi} \overline{f(e^{i\theta})} g(e^{i\theta}) \frac{d\theta}{2\pi}$ and $\|\cdot\|$ the corresponding norm:

Unitary step U , state $\psi \longrightarrow$ Measure $\mu_\psi \longrightarrow$ Schur function f_ψ

$$g_\psi(z) = z \overline{f_\psi(z)} = \sum_{n=1}^{\infty} a_n^\psi z^n$$

Generating function of
first time return amplitudes a_n^ψ

$$R_\psi = \sum_{n=1}^{\infty} |a_n^\psi|^2 = \|g_\psi\|^2 = \|f_\psi\|^2$$

Total return probability

ψ is recurrent $\Leftrightarrow \|f_\psi\| = 1 \Leftrightarrow |f_\psi| = 1$ a.e. in $\mathbb{T} \Leftrightarrow \mu_\psi$ is singular
(f_ψ is inner)

$$\tau_\psi = \sum_{n=1}^{\infty} n |a_n^\psi|^2 = \frac{1}{i} \lim_{r \rightarrow 1} \langle g_\psi(re^{i\theta}), \partial_\theta g_\psi(re^{i\theta}) \rangle = \# \text{supp } \mu_\psi$$

Expected
return time

CONSEQUENCES

- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

CONSEQUENCES

- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

$$\tau_\psi = \# \text{supp } \mu_\psi$$

CONSEQUENCES

- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

$$\tau_\psi = \#\text{supp } \mu_\psi$$

- $\tau_\psi < \infty$ iff μ_ψ is a finite collection of mass points. Hence, the **states with a finite expected return time** are those **spanned by finitely many eigenvectors** of U . In particular, a continuous spectrum for U means that no state can return in a finite expected time.

CONSEQUENCES

- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

$$\tau_\psi = \# \text{supp } \mu_\psi$$

- $\tau_\psi < \infty$ iff μ_ψ is a finite collection of mass points. Hence, the **states with a finite expected return time** are those **spanned by finitely many eigenvectors** of U . In particular, a continuous spectrum for U means that no state can return in a finite expected time.
- In contrast to the classical case, the **expected return time is quantized**. Despite the epithet “*quantum*” of these walks, this is totally unexpected *a priori* because time does not come from an eigenvalue problem in Quantum Mechanics.

CONSEQUENCES

- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

$$\tau_\psi = \#\text{supp } \mu_\psi$$

- $\tau_\psi < \infty$ iff μ_ψ is a finite collection of mass points. Hence, the **states with a finite expected return time** are those **spanned by finitely many eigenvectors** of U . In particular, a continuous spectrum for U means that no state can return in a finite expected time.
- In contrast to the classical case, the **expected return time is quantized**. Despite the epithet “*quantum*” of these walks, this is totally unexpected *a priori* because time does not come from an eigenvalue problem in Quantum Mechanics.
- The above quantization has a **topological reason**: $\tau_\psi < \infty$ iff g_ψ is inner, and then τ_ψ becomes the **winding number** of the boundary values $g_\psi(e^{i\theta}): \mathbb{T} \rightarrow \mathbb{T}$.

CONSEQUENCES

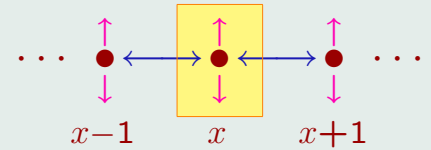
- The **recurrent states** constitute the **singular subspace** of the unitary step U . Thus, an a.c. spectrum for U is equivalent to the absence of recurrent states.

$$\tau_\psi = \# \text{supp } \mu_\psi$$

- $\tau_\psi < \infty$ iff μ_ψ is a finite collection of mass points. Hence, the **states with a finite expected return time** are those **spanned by finitely many eigenvectors** of U . In particular, a continuous spectrum for U means that no state can return in a finite expected time.
- In contrast to the classical case, the **expected return time is quantized**. Despite the epithet “*quantum*” of these walks, this is totally unexpected *a priori* because time does not come from an eigenvalue problem in Quantum Mechanics.
- The above quantization has a **topological reason**: $\tau_\psi < \infty$ iff g_ψ is inner, and then τ_ψ becomes the **winding number** of the boundary values $g_\psi(e^{i\theta}): \mathbb{T} \rightarrow \mathbb{T}$.
- **Quantum recurrence paradox 1**
First time return probabilities can be higher than return probabilities !!!
Nothing prevents $|a_n^\psi|^2 > |\mu_n^\psi|^2$. Indeed, there exist QW with a state ψ such that $a_n^\psi \neq 0$ for all n , but $\mu_n^\psi = 0$ for $n \geq 3$.

Quantum Walks: Subspace recurrence

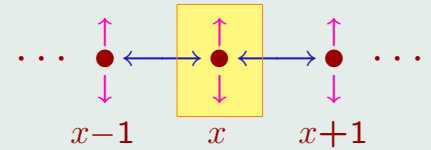
Example Given a state $\psi = \alpha|x, \uparrow\rangle + \beta|x, \downarrow\rangle$ of a coined walk we can be interested in the return probability to the whole site x (instead of to the single state ψ), i.e. the return probability to the subspace $\text{span}\{|x, \uparrow\rangle, |x, \downarrow\rangle\}$



These kind of questions, so natural as the previous ones, lead to the notion of *subspace recurrence* as a generalization of the *state recurrence* already analyzed.

Quantum Walks: Subspace recurrence

Example Given a state $\psi = \alpha|x, \uparrow\rangle + \beta|x, \downarrow\rangle$ of a coined walk we can be interested in the return probability to the whole site x (instead of to the single state ψ), i.e. the return probability to the subspace $\text{span}\{|x, \uparrow\rangle, |x, \downarrow\rangle\}$



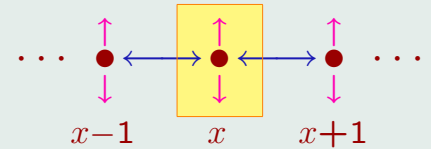
These kind of questions, so natural as the previous ones, lead to the notion of *subspace recurrence* as a generalization of the *state recurrence* already analyzed.

Now we start with a unitary step U and a finite-dimensional subspace V of \mathcal{H} . Given any initial state $\psi \in V$, we search for its return probability to V .

State recurrence corresponds to the particular case $\dim V = 1$.

Quantum Walks: Subspace recurrence

Example Given a state $\psi = \alpha|x, \uparrow\rangle + \beta|x, \downarrow\rangle$ of a coined walk we can be interested in the return probability to the whole site x (instead of to the single state ψ), i.e. the return probability to the subspace $\text{span}\{|x, \uparrow\rangle, |x, \downarrow\rangle\}$



These kind of questions, so natural as the previous ones, lead to the notion of *subspace recurrence* as a generalization of the *state recurrence* already analyzed.

Now we start with a unitary step U and a finite-dimensional subspace V of \mathcal{H} . Given any initial state $\psi \in V$, we search for its return probability to V .

State recurrence corresponds to the particular case $\dim V = 1$.

A precise expression of the **total return probability to V** can be obtained following an approach similar to the case $\dim V = 1$. This simply requires the substitutions

$$\text{Projector onto } \psi \longrightarrow P = \text{Projector onto } V$$

$$\text{Projector onto } \psi^\perp \longrightarrow Q = \text{Projector onto } V^\perp$$

which account for the possible collapses when checking the return to V (instead of ψ).

Quantum Walks: Subspace recurrence

Remind that for state recurrence

- Return amplitude in n steps

$$\mu_n^\psi = \langle \psi | U^n \psi \rangle$$

\rightsquigarrow

$$\text{Prob}(\psi \xrightarrow{n \text{ steps}} \psi) = |\mu_n^\psi|^2$$

- **First time** return amplitude in n steps

$$a_n^\psi = \langle \psi | U(QU)^{n-1} \psi \rangle$$

\rightsquigarrow

$$\text{Prob}(\psi \xrightarrow[n \text{ steps}]{1^{\text{st}} \text{ time}} \psi) = |a_n^\psi|^2$$

Quantum Walks: Subspace recurrence

In the case of subspace recurrence

- **Return amplitude in n steps** $\mu_n^V = PU^n P \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} V) = \|\mu_n^V \psi\|^2$
- **First time return amplitude in n steps** $a_n^V = PU(QU)^{n-1} P \rightsquigarrow \text{Prob}(\psi \xrightarrow[1^{\text{st}} \text{ time}]{n \text{ steps}} V) = \|a_n^V \psi\|^2$

Quantum Walks: Subspace recurrence

In the case of subspace recurrence

- **Return amplitude in n steps** $\mu_n^V = PU^n P \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} V) = \|\mu_n^V \psi\|^2$
- **First time return amplitude in n steps** $a_n^V = PU(QU)^{n-1}P \rightsquigarrow \text{Prob}(\psi \xrightarrow[1^{\text{st}} \text{ time}]{n \text{ steps}} V) = \|a_n^V \psi\|^2$

Now μ_n^V, a_n^V are operators on V that we identify with matrices choosing a basis of V

Quantum Walks: Subspace recurrence

In the case of subspace recurrence

- **Return amplitude in n steps** $\mu_n^V = PU^n P \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} V) = \|\mu_n^V \psi\|^2$

- **First time return amplitude in n steps** $a_n^V = PU(QU)^{n-1}P \rightsquigarrow \text{Prob}(\psi \xrightarrow[n \text{ 1st time}]{n \text{ steps}} V) = \|a_n^V \psi\|^2$

Now μ_n^V, a_n^V are operators on V that we identify with matrices choosing a basis of V

$$R_V(\psi) = \sum_{n=1}^{\infty} \|a_n^V \psi\|^2 \quad \text{Total } V\text{-return probability}$$

$$\psi \text{ is } V\text{-recurrent} \Leftrightarrow R_V(\psi) = 1$$

$$\tau_V(\psi) = \sum_{n=1}^{\infty} n \|a_n^V \psi\|^2 \quad \text{Expected } V\text{-return time}$$

Quantum Walks: Subspace recurrence

In the case of subspace recurrence

- **Return amplitude in n steps** $\mu_n^V = PU^nP \rightsquigarrow \text{Prob}(\psi \xrightarrow{n \text{ steps}} V) = \|\mu_n^V \psi\|^2$
- **First time return amplitude in n steps** $a_n^V = PU(QU)^{n-1}P \rightsquigarrow \text{Prob}(\psi \xrightarrow[1^{\text{st time}}]{n \text{ steps}} V) = \|a_n^V \psi\|^2$

Now μ_n^V, a_n^V are operators on V that we identify with matrices choosing a basis of V

$$R_V(\psi) = \sum_{n=1}^{\infty} \|a_n^V \psi\|^2 \quad \text{Total } V\text{-return probability}$$

$$\psi \text{ is } V\text{-recurrent} \Leftrightarrow R_V(\psi) = 1$$

$$\tau_V(\psi) = \sum_{n=1}^{\infty} n \|a_n^V \psi\|^2 \quad \text{Expected } V\text{-return time}$$

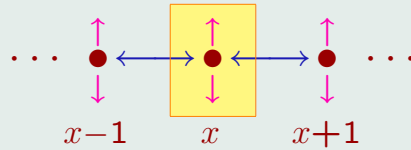
We can rewrite $R_V(\psi) = \langle \psi | \mathbf{R}_V \psi \rangle$ as a quadratic form in V giving by

$$\mathbf{R}_V = \sum_{n=1}^{\infty} (a_n^V)^\dagger a_n^V$$

Example: Site recurrence in a translation invariant coined walk

Consider the site subspace $V_x = \text{span}\{|x, \uparrow\rangle, |x, \downarrow\rangle\}$ in the previous 1D model

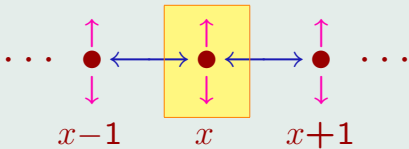
$$X = \mathbb{Z} \times \{\uparrow, \downarrow\}$$



$$C = \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix} \text{ unitary coin}$$

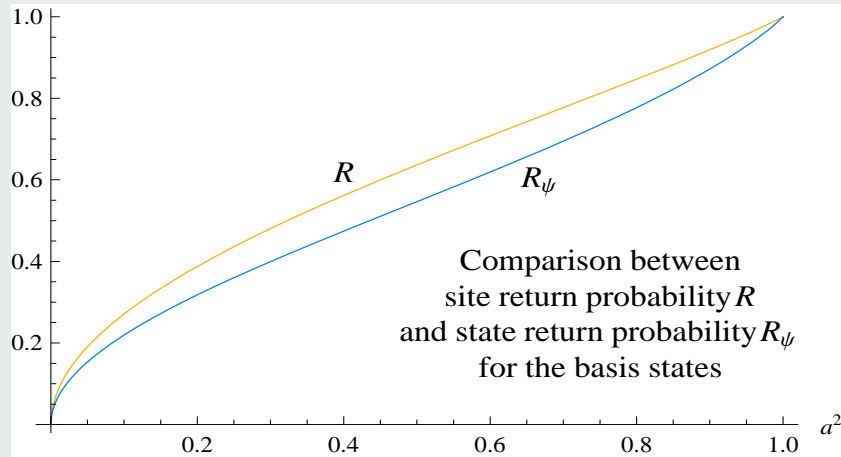
Example: Site recurrence in a translation invariant coined walk

Consider the site subspace $V_x = \text{span}\{|x, \uparrow\rangle, |x, \downarrow\rangle\}$ in the previous 1D model

$$X = \mathbb{Z} \times \{\uparrow, \downarrow\}$$


$$C = \begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix} \text{ unitary coin}$$

$$R_{V_x} = \begin{pmatrix} R & 0 \\ 0 & R \end{pmatrix} \text{ is a scalar matrix with } R = \frac{\rho a + (1 - 2\rho^2) \arcsin a}{\frac{\pi}{2} a^2} \quad \begin{cases} a = |c_{12}| \\ \rho = \sqrt{1 - a^2} \end{cases}$$



Comparison between state and site recurrence for $\psi = |x, \uparrow\rangle, |x, \downarrow\rangle$
 Note that $R_\psi \leq R = R_{V_x}(\psi)$ as one should naively expect

Quantum Walks and **matrix** Schur functions

The spectral characterization of subspace recurrence requires the association of a matrix measure and a matrix Schur function with each subspace.

Quantum Walks and **matrix** Schur functions

The spectral characterization of subspace recurrence requires the association of a matrix measure and a matrix Schur function with each subspace.

As in the case of a single state, the unitarity of U ensures for any **subspace** $V \subset \mathcal{H}$ the existence of a **spectral matrix measure** μ_V on \mathbb{T} such that

n -th moment $\int_{\mathbb{T}} t^n d\mu_V(t) = PU^nP = \mu_n^V$ return (matrix) amplitude to V in n steps

Quantum Walks and **matrix** Schur functions

The spectral characterization of subspace recurrence requires the association of a matrix measure and a matrix Schur function with each subspace.

As in the case of a single state, the unitarity of U ensures for any subspace $V \subset \mathcal{H}$ the existence of a **spectral matrix measure** μ_V on \mathbb{T} such that

n -th moment $\int_{\mathbb{T}} t^n d\mu_V(t) = PU^nP = \mu_n^V$ return (matrix) amplitude to V in n steps

This allows to associate with V a **matrix Schur function** f_V

$$f_V(z) = z^{-1}(F_V(z) - I)(F_V(z) + I)^{-1} \quad F_V(z) = \int_{\mathbb{T}} \frac{t+z}{t-z} d\mu_V(t)$$

which is analytic on the unit disk \mathbb{D} and satisfies $\|f_V\| \leq 1$ in \mathbb{D} .

Quantum Walks and **matrix** Schur functions

The spectral characterization of subspace recurrence requires the association of a matrix measure and a matrix Schur function with each subspace.

As in the case of a single state, the unitarity of U ensures for any **subspace** $V \subset \mathcal{H}$ the existence of a **spectral matrix measure** μ_V on \mathbb{T} such that

n -th moment $\int_{\mathbb{T}} t^n d\mu_V(t) = PU^nP = \mu_n^V$ return (matrix) amplitude to V in n steps

This allows to associate with V a **matrix Schur function** f_V

$$f_V(z) = z^{-1}(F_V(z) - I)(F_V(z) + I)^{-1} \quad F_V(z) = \int_{\mathbb{T}} \frac{t+z}{t-z} d\mu_V(t)$$

which is analytic on the unit disk \mathbb{D} and satisfies $\|f_V\| \leq 1$ in \mathbb{D} .

The theory of matrix Schur functions, linked to the matrix version of OP on \mathbb{T} , has experienced a strong development influenced by the needs of electrical engineering, signal transmission and processing, prediction theory for stochastic processes, ...

A list to which we should add from now the issue of recurrence in QW:

Matrix Schur functions are the math objects which best codify **subspace recurrence**.

Quantum Walks: Spectral characterization of **subspace** recurrence

Denoting $\langle\langle f, g \rangle\rangle = \int_0^{2\pi} f(e^{i\theta})^\dagger g(e^{i\theta}) \frac{d\theta}{2\pi}$ and $\|\cdot\|$ the corresponding matrix “norm”:

Unitary step U , subspace $V \longrightarrow$ Matrix measure $\mu_V \longrightarrow$ Matrix Schur function f_V

$$g_V(z) = z f_V^\dagger(z) = \sum_{n=1}^{\infty} a_n^V z^n$$

Generating function of **first time**
return matrix amplitudes a_n^V

$$R_V = \|\|g_V\|\|^2 = \|\|f_V^\dagger\|\|^2 \rightsquigarrow R_V(\psi) = \langle\psi|R_V\psi\rangle$$

**Total V -return
probability**

**Any $\psi \in V$ is
 V -recurrent** $\Leftrightarrow \|\|f_V\|\| = 1 \Leftrightarrow f_V$ is unitary a.e. in $\mathbb{T} \Leftrightarrow \mu_V$ is singular
(f_ψ is inner)

$$\tau_V^r = \frac{1}{i} \langle\langle g_V(re^{i\theta}), \partial_\theta g_V(re^{i\theta}) \rangle\rangle \rightsquigarrow \tau_V(\psi) = \lim_{r \rightarrow 1} \langle\psi|\tau_V^r\psi\rangle$$

**Expected
 V -return time**

$\tau_V(\psi) < \infty$ for all $\psi \in V \Leftrightarrow \mu_V$ is a finite combination of mass points

Quantum Walks: State recurrence \longleftrightarrow Subspace recurrence

Despite the previous analogies, there are many differences and unexpected relations

- Recurrent states need a singular subspace for U
- NO, unless V -recurrent states cover all of V

Quantum Walks: State recurrence \longleftrightarrow Subspace recurrence

Despite the previous analogies, there are many differences and unexpected relations

- Recurrent states need a singular subspace for U
- NO, unless V -recurrent states cover all of V

- States returning in a finite expected time need eigenvectors of U
- NO, unless the states returning to V in a finite expected time cover all of V

Quantum Walks: State recurrence \longleftrightarrow Subspace recurrence

Despite the previous analogies, there are many differences and unexpected relations

- Recurrent states need a singular subspace for U
- NO, unless V -recurrent states cover all of V

- States returning in a finite expected time need eigenvectors of U
- NO, unless the states returning to V in a finite expected time cover all of V

When the the expected return time is finite:

- $\tau_\psi =$ **winding number of** $g_\psi(e^{i\theta})$ is an INTEGER TOPOLOGICAL INVARIANT
- NO. However,

Quantum Walks: State recurrence \longleftrightarrow Subspace recurrence

Despite the previous analogies, there are many differences and unexpected relations

- Recurrent states need a singular subspace for U
- NO, unless V -recurrent states cover all of V

- States returning in a finite expected time need eigenvectors of U
- NO, unless the states returning to V in a finite expected time cover all of V

When the the expected return time is finite:

- $\tau_\psi =$ **winding number of** $g_\psi(e^{i\theta})$ is an INTEGER TOPOLOGICAL INVARIANT
- NO. However,
 - ▶ $\tau_V(\psi)$ is a GEOMETRICAL INVARIANT of the curve $g_V(e^{i\theta})\psi$ which reflects the non-trivial geometry of the projective space of states (**Berry's phase**)

Quantum Walks: State recurrence \longleftrightarrow Subspace recurrence

Despite the previous analogies, there are many differences and unexpected relations

- Recurrent states need a singular subspace for U
- NO, unless V -recurrent states cover all of V

- States returning in a finite expected time need eigenvectors of U
- NO, unless the states returning to V in a finite expected time cover all of V

When the the expected return time is finite:

- $\tau_\psi =$ **winding number of** $g_\psi(e^{i\theta})$ is an INTEGER TOPOLOGICAL INVARIANT
- NO. However,
 - ▶ $\tau_V(\psi)$ is a GEOMETRICAL INVARIANT of the curve $g_V(e^{i\theta})\psi$ which reflects the non-trivial geometry of the projective space of states (**Berry's phase**)
 - ▶ The V -AVERAGE of $\tau_V(\psi)$ is linked to a true TOPOLOGICAL INVARIANT

$$\int_V \tau_V(\psi) d\psi = \frac{\text{winding number of } \det g_V(e^{i\theta})}{\dim V} \quad (\text{RATIONAL NUMBER})$$

- **Quantum recurrence paradox 2**

State return probabilities can be higher than subspace return probabilities !!!



The above figure shows an example of this kind of phenomenon in the case of a 1D coined walk.

It represents as a function of t the state and site return probability of the states $\psi(t) = \cos t|x, \uparrow\rangle + \sin t|x, \downarrow\rangle$ lying in the same site x .

We can see that the state return probability (in blue) is occasionally bigger than the site return probability (in red).

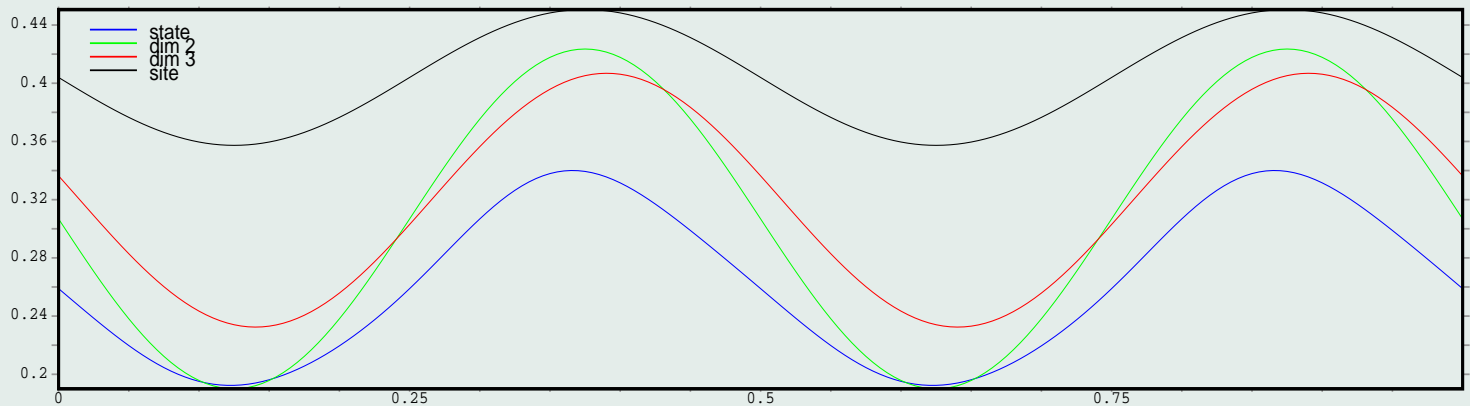
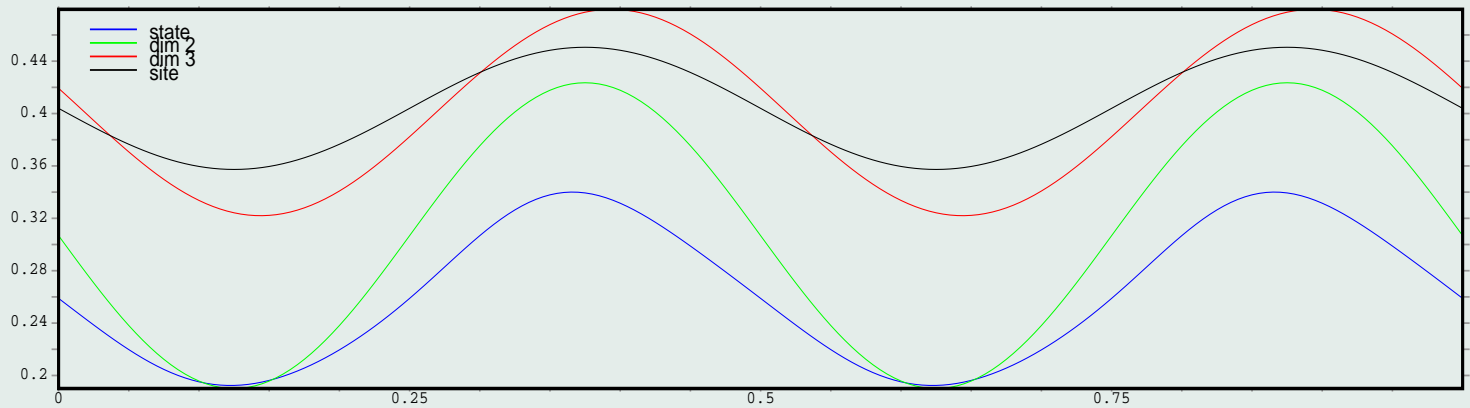


Figure 1: Similar figures for a 2D QW in a square lattice. They compare, for a certain curve of states $\psi(t)$, the return probability to some nested subspaces of dimension 1 (the state), 2, 3, and 4 (the site). The return probability does not necessarily increase when enlarging the subspace.

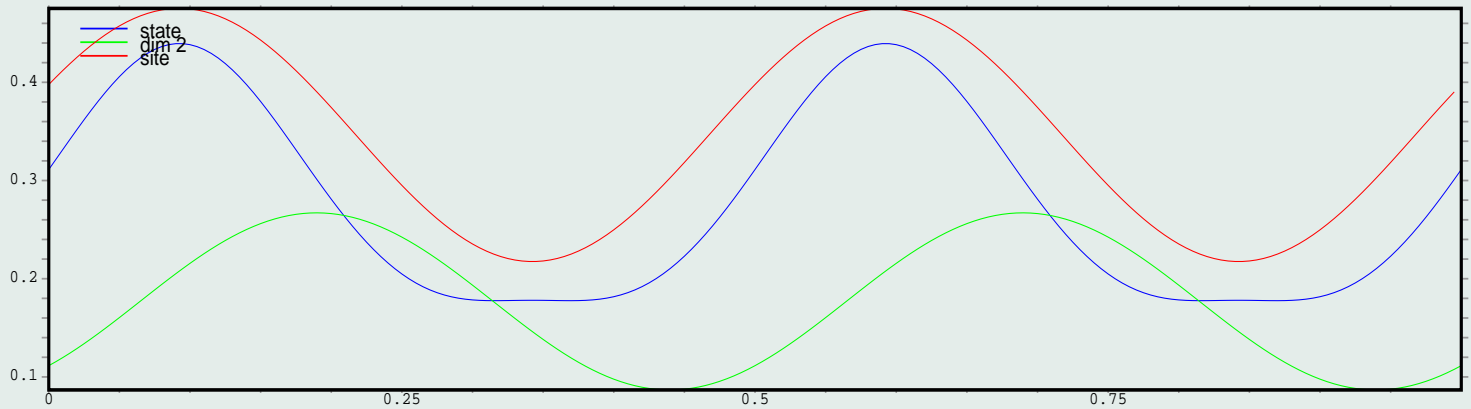
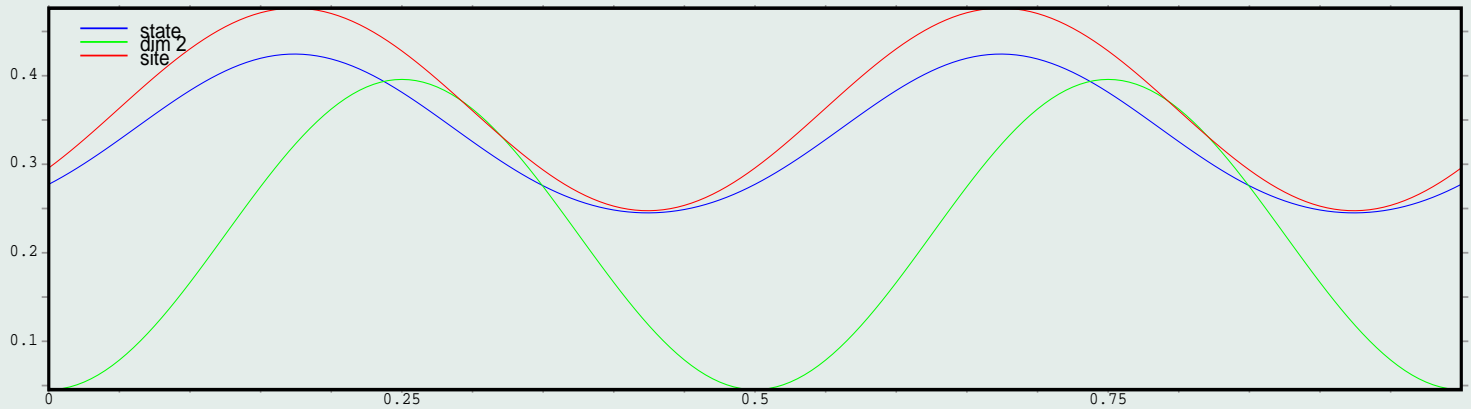


Figure 2: The QW lives now in a 2D hexagonal lattice. The figures compare the return probability of a curve of states to some nested subspaces of dimension 1 (the state), 2 and 3 (the site). The relation between the cases of dimension 1 and 2 is dramatic because it is most of the times the opposite of what one should naively expect.

Quantum Walks: Comments on recurrence

T_ψ INTEGER or INFINITE

Quantum recurrence paradoxes

↪ Experimental validation?

Quantum Walks: Comments on recurrence

τ_ψ INTEGER or INFINITE

Quantum recurrence paradoxes

↪ Experimental validation?

$$\int_V \tau_V(\psi) d\psi \text{ RATIONAL or INFINITE is equivalent to}$$
$$\sum_{n=1}^{\infty} n \operatorname{Tr}[PU(QU)^{n-1}P] \text{ INTEGER or INFINITE}$$

for any unitary U and orthogonal projectors $P, Q = I - P$

↪ New result in Operator Theory?

Quantum Walks: Comments on recurrence

τ_ψ INTEGER or INFINITE

Quantum recurrence paradoxes

↪ Experimental validation?

$$\int_V \tau_V(\psi) d\psi \quad \text{RATIONAL or INFINITE} \quad \text{is equivalent to}$$
$$\sum_{n=1}^{\infty} n \operatorname{Tr}[PU(QU)^{n-1}P] \quad \text{INTEGER or INFINITE}$$

for any unitary U and orthogonal projectors $P, Q = I - P$

↪ New result in Operator Theory?

Dirac equation

↪ Continuous time version?

Quantum Walks: Returns to Schur and OP theory

Schur / OP \longrightarrow QW

Quantum Walks: Returns to Schur and OP theory

Schur / OP \longrightarrow QW

Schur / OP \longleftarrow ? QW

Quantum Walks: Returns to Schur and OP theory

Schur / OP \longrightarrow QW

Schur / OP \longleftarrow ? QW

KHRUSHCHEV \longleftarrow FEYNMAN
formulas diagrams

Quantum Walks: Returns to Schur and OP theory

Schur / OP \longrightarrow QW

Schur / OP \longleftarrow ? QW

KHRUSHCHEV \longleftarrow FEYNMAN
formulas diagrams

ILAS Meeting, Providence, RI, USA

June 3-7, 2013